

**Joint seminar on cosmology and astroparticle physics  
Tõravere, 26-27 February 2007**

# **Topics in String Cosmology**

**(How to get accelerated expansion)**

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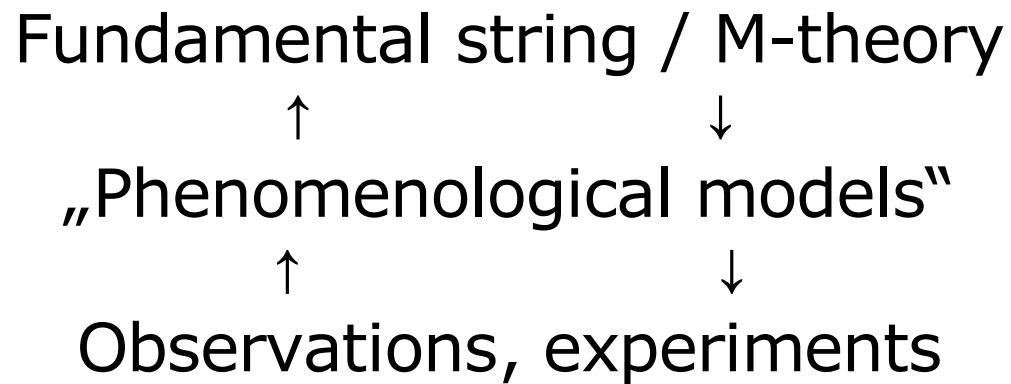


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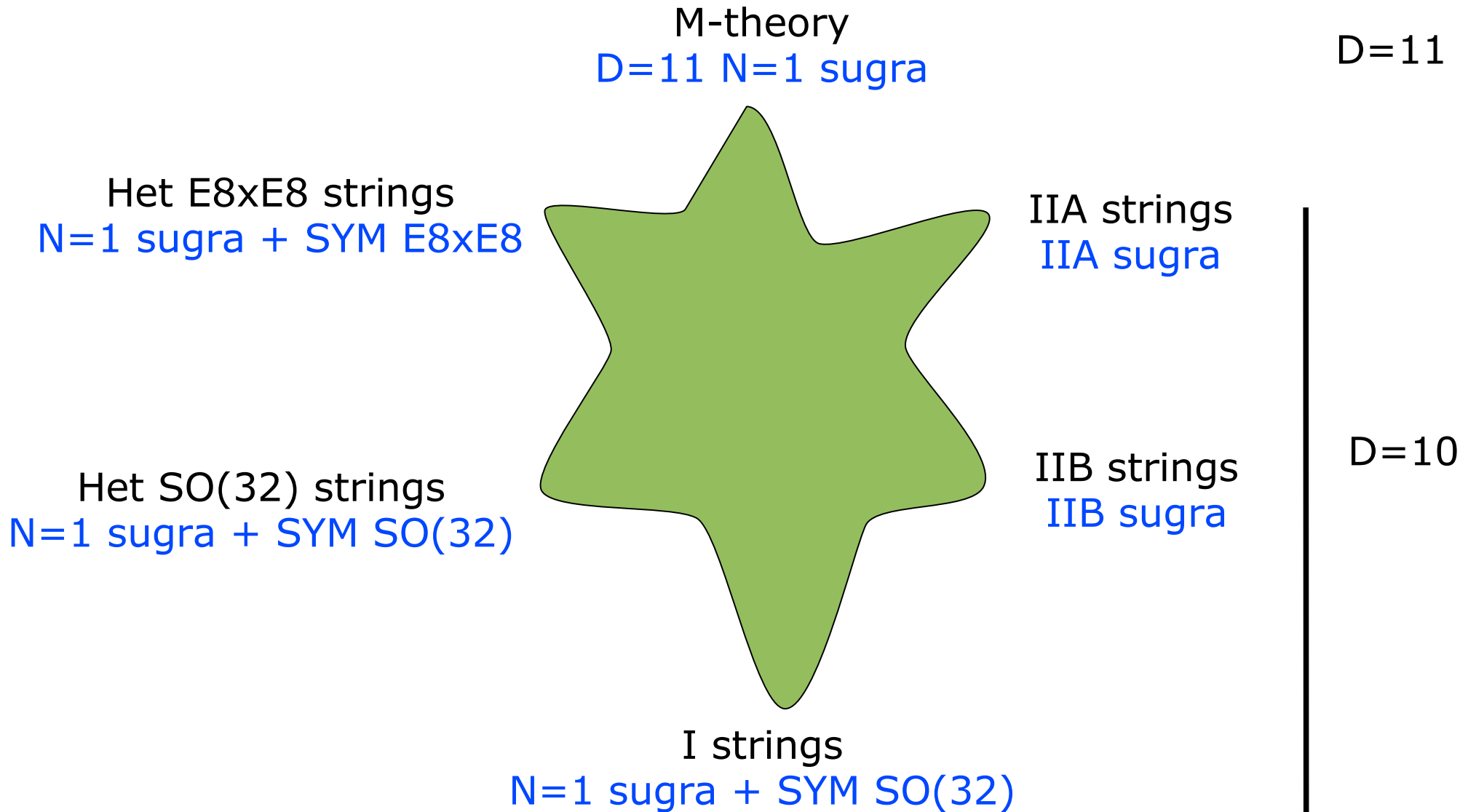
<http://hexagon.fi.tartu.ee/>

## **Standard cosmology**, accelerated expansion

- *Ad hoc*: Inflaton potential, value of cosmological constant



# Web of String Theories



## Type II String Theory

$$S_{NS-NS} = \frac{1}{\kappa_{10}^2} \int d^{10}x \sqrt{-g} e^{-2\Phi} \left( R + 4(\partial\Phi)^2 - \frac{1}{2}|H_3|^2 \right)$$

$$S_{R-R} = -\frac{1}{4\kappa_{10}^2} \int d^{10}x \begin{cases} |F_2|^2 + |F_4|^2, & IIA \\ |F_1|^2 + |F_3|^2 + \frac{1}{2}|F_5|^2, & IIB \end{cases} \quad F_6 = *F_4, \quad F_8 = *F_2, \quad F_5 = *F_5$$

- Dp-branes: dynamical (p + 1)-dimensional objects (with p spatial dimensions)
- $F_n$  couples to the Dp-brane with  $p = n-2$

### To obtain a 4D observable world

- Compactification
- Braneworld large extra dimensions

## **No-go theorem**

Compactifications of 10, 11 dim SUGRA on

- Time-independent
- Compact
- Smooth

internal space do not yield de Sitter vacua.

(Gibbons 1987, Maldacena, NuneZ 2001)

## **How to derive a scalar potential that leads to accelerated expansion?**

- Potential induced by compactification
  - With fluxes
  - Time-dependent hyperbolic manifold
  - Time-dependent singular manifold
- Brane-antibrane potential

## A simple model

$$S^{(d+m)} = \frac{1}{\kappa^{(d+m)}} \int d^d x d^m y \sqrt{-g} \left( R - \frac{1}{2m!} (F_m)^2 \right) \quad F_m = b \text{Vol}_{\text{int}}$$

Compactify on a maximally symmetric space

$$ds^2 = \tilde{g}_{\mu\nu}(x) dx^\mu dx^\nu + \hat{g}_{mn}(x, y) dy^m dy^n \quad \hat{g}_{mn} = e^{2\beta(x)} \hat{\omega}_{mn}(y)$$

Space-time macroscopic d-dim	Internal space compact m-dim
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$$\hat{R}_{mn}(\hat{\omega}) = k_1 (m-1) \hat{\omega}_{mn}(y)$$

$$\int d^m y \sqrt{-\hat{\omega}(y)} = 1$$

Effective action in 4D

$$S^{(4)} = \int d^4 x \sqrt{-g} \left( \frac{1}{2} R - (\partial\phi)^2 - 2V(\phi) \right) \quad V(\phi) = -k_1 e^{-2c\phi} + \frac{\tilde{b}^2}{2} e^{-\frac{6}{c}\phi}$$

$$c = \sqrt{\frac{m+2}{m}}, \quad b^2 =: 4\tilde{b}^2 \left( \frac{4}{m(m-1)} \right)^{-3/c^2}, \quad \beta(x) = \frac{2}{mc} \phi(x) + \frac{1}{m+2} \ln \left( \frac{m(m-1)}{4} \right)$$

## **FRW cosmology of the model**

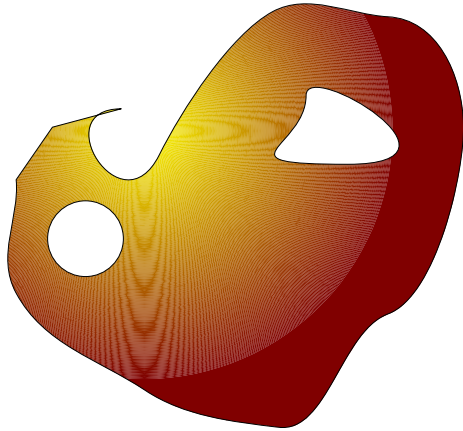
- Acceleration: transient, eternal or eternally cyclic going to zero asymptotically (not good enough for inflation)
- Runaway potential
- Modulus not stabilized, spontaneous decompactification

(L. Järv, T. Mohaupt, F. Saueressig 2004)

## Moduli trapping

- Fluxes
- Non-perturbative effects  
(S. Kachru, R. Kallosh, A. Linde, S. Trivedi 2003)
- ESP's (extra species points, enhanced symmetry points)  
(M. Brändle, A. Lukas 2002,  
L. Järv, T. Mohaupt, F. Saueressig 2003,  
L. Kofman, A. Linde, X. Liu, A. Maloney, L. McAllister, E. Silverstein  
2004)

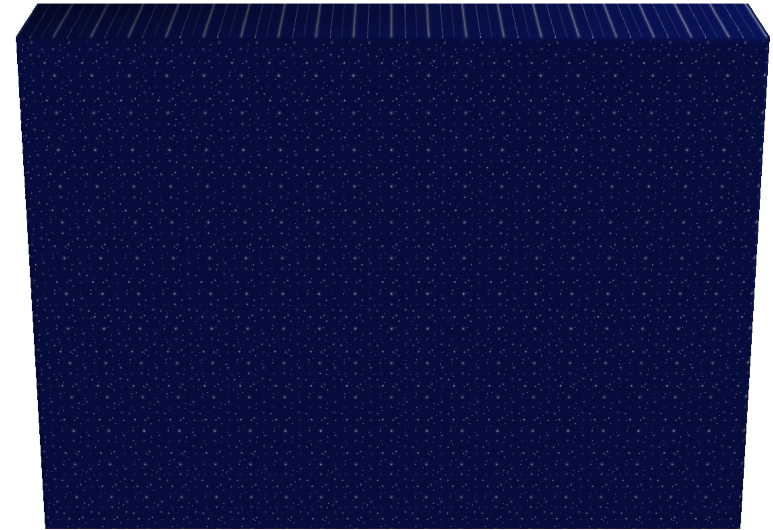
Internal space (microscopic)



Moduli –  
parameters describing the  
internal geometry and  
topology

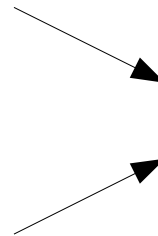
Fluxes in the internal space

Spacetime (macroscopic)



Scalar fields in spacetime  
(with a potential)

Cosmological constant –  
minimum of the potential



## String theory landscape

Live in a local positive minimum

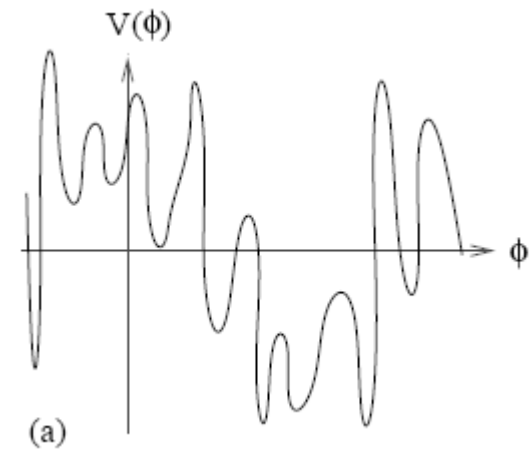
- Metastable, but long lifetime
- Eternal inflation

Number of different „vacua“

- e.g. in Calabi-Yau compactifications of type IIB string theory:  $10^{500}$   
(M. Douglas 2003)

What to do?

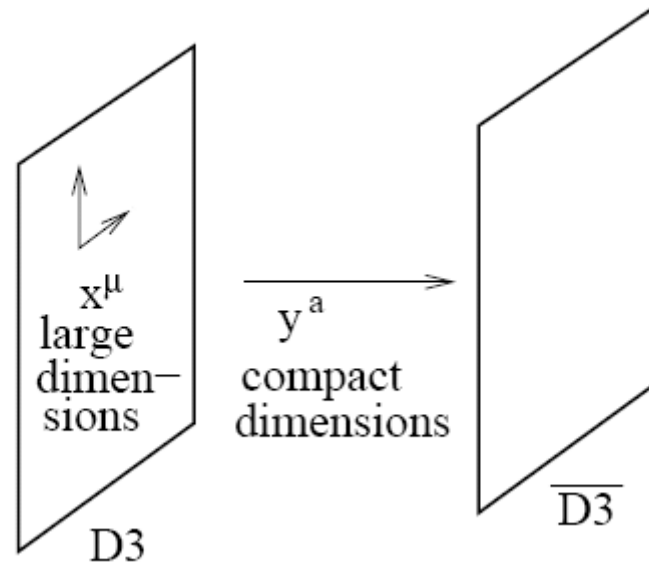
- Statistics, anthropic selection
- Dynamical mechanism?



## Brane and antibrane

$$\tau_p = \frac{M_s^{p+1}}{(2\pi)^p g_s}$$

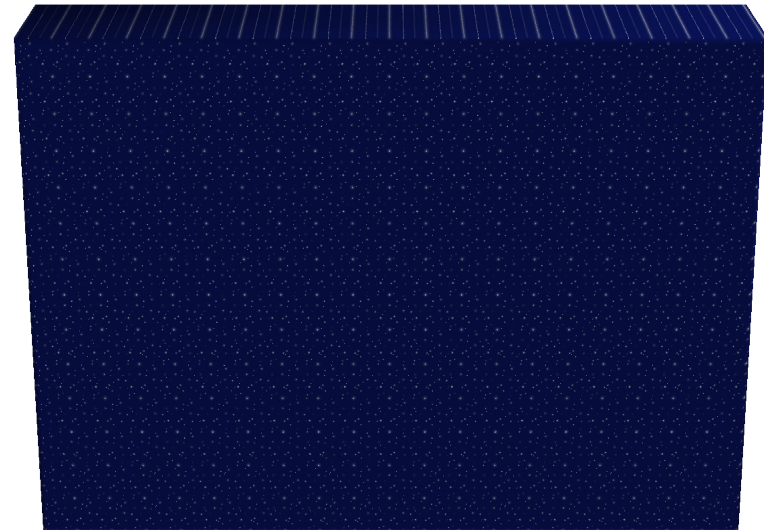
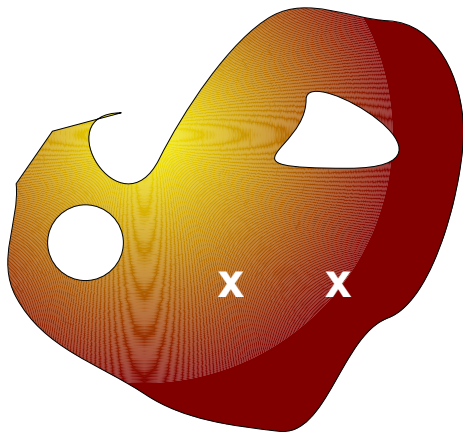
$$\mu_p = g_s \tau_p$$



$$V_{\text{tot}} = \frac{G_{10}}{\frac{1}{4}\pi^2 r^4} \left( -\tau_3^2 \pm \frac{\mu_3^2}{g_s^2} \right) \begin{cases} \text{upper sign for D3-D3} \\ \text{lower sign for D3-}\overline{\text{D3}} \end{cases}$$

$$G_{10} = (2\pi)^6 g_s^2 / (8M_s^8)$$

## Brane inflation



D3 brane and antibrane

Separated in internal space, but

fill the 4D spacetime

For small  
velocities  
large  
separation

$$\phi = \sqrt{\tau_3} |\vec{y} - \vec{\bar{y}}|$$

$$c = \frac{4}{\pi^2} G_{10} \tau_3^4$$

$$\mathcal{L} = \frac{1}{2} \dot{\phi}^2 + 2 \left( \tau_3 - \frac{c}{\phi^4} \right)$$

At close separation brane and antibrane annihilate

Unfortunately not slow roll

## Various models

- Warped brane-antibrane inflation

S. Kachru, R. Kallosh, A. Linde, J. M. Maldacena, L. McAllister, S. P. Trivedi 2003

- Superpotential corrections
- Tuning the length of the throat
- Multibrane inflation

- DBI inflation (D-acceleration), fast roll

E. Silverstein, D. Tong, M. Alishahiha 2004

- Shift symmetry, D3-D7 inflation

K. Dasgupta, C. Herdeiro, S. Hirano and R. Kallosh 2002

J. P. Hsu and R. Kallosh 2004

- Racetrack and Kähler moduli inflation

J. J. Blanco-Pillado 2004